

# APMA 2560 Homework #2 Solutions

February 15, 2008

1.) We have shown in class that if the operator  $\mathcal{L}$  and the projection  $\mathcal{P}_N$  commute as

$$\mathcal{P}_N \mathcal{L} \mathcal{P}_N u(x) = \mathcal{P}_N \mathcal{L} u(x),$$

then there is no truncation error and one can solve the problem exactly as long as the initial conditions are exactly represented within  $B_N$ .

For the Fourier expansion

$$\mathcal{P}_N u(x) = \sum_{|n| \leq N/2} \hat{u}_n e^{inx}, \quad \hat{u}_n = \frac{1}{2\pi} \int_0^{2\pi} u(x) e^{-inx} dx,$$

and  $u(x)$  being periodic on  $x \in [0, 2\pi]$ .

a) What is  $B_N$  in this case?

b) Prove that if the operator is

$$\mathcal{L} = \frac{d^q}{dx^q},$$

the truncation error vanishes.

c) Prove that if the operator is

$$\mathcal{L} = \sin(x) \frac{d}{dx},$$

the truncation error does not vanish.

## Solution:

a.) In the statement of the problem, we assume that  $N$  is even. In this case, the space  $B_N$  is given by,

$$B_N = \text{span}\{e^{inx} : |n| \leq N/2\},$$

where the linear span is taken over the field of complex numbers.

b.) We take

$$\mathcal{L} = \frac{d^q}{dx^q},$$

and assume  $u(x)$  any  $L^2([0, 2\pi])$  function. Then  $u(x)$  admits a Fourier expansion:

$$u(x) = \sum_{|n| < \infty} \hat{u}_n e^{inx}. \tag{1}$$

Given this expansion, then we have

$$\mathcal{P}_N u(x) = \sum_{|n| \leq N/2} \hat{u}_n e^{inx}. \quad (2)$$

Of course, this implies that for all  $u \in B_N$ , we have  $\mathcal{P}_N u = u$ . We also have that for any  $u \in L^2$ ,

$$\mathcal{L} u(x) = \sum_{|n| < \infty} (in)^q \hat{u}_n e^{inx}.$$

Armed with these two statements, we can explicitly compute

$$\begin{aligned} \mathcal{P}_N \mathcal{L} \mathcal{P}_N u(x) &= \mathcal{P}_N \mathcal{L} \sum_{|n| \leq N/2} \hat{u}_n e^{inx} \\ &= \mathcal{P}_N \sum_{|n| \leq N/2} (in)^q \hat{u}_n e^{inx} \\ &= \sum_{|n| \leq N/2} (in)^q \hat{u}_n e^{inx}. \end{aligned}$$

We also compute

$$\begin{aligned} \mathcal{P}_N \mathcal{L} u(x) &= \mathcal{P}_N \sum_{|n| < \infty} (in)^q \hat{u}_n e^{inx} \\ &= \sum_{|n| \leq N/2} (in)^q \hat{u}_n e^{inx} \end{aligned}$$

We can now explicitly see that

$$\mathcal{P}_N \mathcal{L} \mathcal{P}_N u(x) = \mathcal{P}_N \mathcal{L} u(x),$$

which shows that the truncation error is zero. In this case, the space  $B_N$  is closed under the operator  $\mathcal{L}$ .

c.) We now take

$$\mathcal{L} = \sin(x) \frac{d}{dx}.$$

We still assume an expansion of  $u$  as given in (1) and the projection result (2) also still holds. For any  $u \in L^2([0, 2\pi])$ , we can compute

$$\begin{aligned} \mathcal{L} u(x) &= \sin(x) \sum_{|n| < \infty} (in) \hat{u}_n e^{inx} \\ &= \sum_{|n| < \infty} \frac{n}{2} \hat{u}_n (e^{ix} - e^{-ix}) e^{inx} \\ &= \sum_{|n| < \infty} \left( \frac{n-1}{2} \hat{u}_{n-1} - \frac{n+1}{2} \hat{u}_{n+1} \right) e^{inx}. \end{aligned}$$

Note that this already shows that the space  $B_N$  is not closed under the operator  $\mathcal{L}$ . In any case,

we compute

$$\begin{aligned}
\mathcal{P}_N \mathcal{L} \mathcal{P}_N u(x) &= \mathcal{P}_N \mathcal{L} \sum_{|n| \leq N/2} \hat{u}_n e^{inx} \\
&= \mathcal{P}_N \left[ \sum_{|n| \leq N/2-1} \left( \frac{n-1}{2} \hat{u}_{n-1} - \frac{n+1}{2} \hat{u}_{n+1} \right) e^{inx} \right. \\
&\quad + \frac{N/2-1}{2} \hat{u}_{N/2-1} e^{ixN/2} - \frac{-N/2+1}{2} \hat{u}_{-N/2+1} e^{-ixN/2} \\
&\quad \left. + \frac{N/2}{2} \hat{u}_{N/2} e^{ix(N+1)/2} - \frac{-N/2}{2} \hat{u}_{-N/2} e^{-ix(N+1)/2} \right] \\
&= \sum_{|n| \leq N/2-1} \left( \frac{n-1}{2} \hat{u}_{n-1} - \frac{n+1}{2} \hat{u}_{n+1} \right) e^{inx} \\
&\quad + \frac{N/2-1}{2} \hat{u}_{N/2-1} e^{ixN/2} - \frac{-N/2+1}{2} \hat{u}_{-N/2+1} e^{-ixN/2}.
\end{aligned}$$

We also have

$$\begin{aligned}
\mathcal{P}_N \mathcal{L} u(x) &= \mathcal{P}_N \sum_{|n| < \infty} \left( \frac{n-1}{2} \hat{u}_{n-1} - \frac{n+1}{2} \hat{u}_{n+1} \right) e^{inx} \\
&= \sum_{|n| \leq N/2} \left( \frac{n-1}{2} \hat{u}_{n-1} - \frac{n+1}{2} \hat{u}_{n+1} \right) e^{inx}.
\end{aligned}$$

Comparing these last two derived expressions, we see that  $\mathcal{P}_N \mathcal{L} \mathcal{P}_N u(x)$  is “missing” two terms: a  $\hat{u}_{N/2+1} e^{ixN/2}$  and a  $\hat{u}_{-N/2-1} e^{-ixN/2}$ . Thus, in this case the truncation error

$$\mathcal{P}_N \mathcal{L}(I - \mathcal{P}_N)u(x)$$

is not identically zero for any  $u$ . Note that if we assume that  $u \in B_{N-1}$ , however, then the truncation error is zero.

2.) For the linear wave equation

$$\frac{\partial u}{\partial t} = -c \frac{\partial u}{\partial x}, \quad x \in [0, 2\pi]$$

we have considered ways of approximating the derivative using a Fourier series. Using the orthogonality of the basis, we have recovered that

$$\frac{d\hat{u}_n}{dt} = -i c n \hat{u}_n, \quad |n| \leq N/2. \quad (3)$$

a) Write this on the generic form

$$\frac{d\mathbf{u}_N}{dt} = \mathcal{L}_N \mathbf{u}_N,$$

and define what  $\mathbf{u}_N$  and  $\mathcal{L}_N$  mean.

b) Prove that the scheme is stable in the sense of

$$\|\exp(\mathcal{L}_N t)\| \leq C(t),$$

and specify what  $C(t)$  is.

**Solution:**

a.) We define the modal coefficients  $\hat{u}_n$  to be

$$\hat{u}_n = \frac{1}{2\pi} \int_0^{2\pi} u(x) e^{-inx} dx.$$

For a given  $N$  even, we arrange them in a vector as follows:

$$\mathbf{u}_N = \begin{pmatrix} \hat{u}_{-N/2} \\ \hat{u}_{-N/2+1} \\ \vdots \\ \hat{u}_{N/2-1} \\ \hat{u}_{N/2} \end{pmatrix}.$$

We define  $\mathcal{L}_N$  to be a diagonal  $(N+1) \times (N+1)$  matrix:

$$\mathcal{L}_N = \begin{pmatrix} \frac{icN}{2} & 0 & 0 & \dots & \dots & 0 \\ 0 & \frac{ic(N-1)}{2} & 0 & \dots & \dots & 0 \\ 0 & 0 & \frac{ic(N-2)}{2} & & & \\ \vdots & \vdots & & \ddots & & \vdots \\ \vdots & \vdots & & & -\frac{ic(N-1)}{2} & 0 \\ 0 & 0 & \dots & \dots & 0 & -\frac{icN}{2} \end{pmatrix}$$

Then the semi-discrete form (3) can be written as

$$\frac{d\mathbf{u}_N}{dt} = \mathcal{L}_N \mathbf{u}_N.$$

b.) For a diagonalizable matrix  $A = S \Lambda S^{-1}$ , one way of defining the matrix exponential is

$$\exp(A) = S \exp(\Lambda) S^{-1},$$

where  $\exp(\Lambda)$ , the exponential of a diagonal matrix is simply the element-by-element exponential of the diagonal entries, and zero elsewhere. Since the operator  $\mathcal{L}_N$  is already diagonal, we can directly apply this. For the norm, the 2-norm of a matrix is defined as

$$\|A\| = \sup_{x \neq 0} \frac{\|Ax\|}{\|x\|},$$

where the norms on the elements  $x$  and  $Ax$  in this case can be taken as Euclidean 2-norms. One can then show from this that for a diagonalizable matrix  $A$  with eigenvalues  $\lambda_i$ , the matrix 2-norm is the same as the spectral radius  $\rho(A)$ :

$$\|A\| = \rho(A) = \max_i |\lambda_i|.$$

In other words, we have that

$$\|\exp(\Lambda t)\| = \max_{|n| \leq N/2} |e^{-icnt}| = 1.$$

So for the ODE (3), we have that

$$\|\exp(\Lambda t)\| = C(t) \equiv 1.$$

3.) Using the Fourier differentiation matrix, compute the derivative of the following functions

(a)  $f(x) = \exp(\cos(4x))$

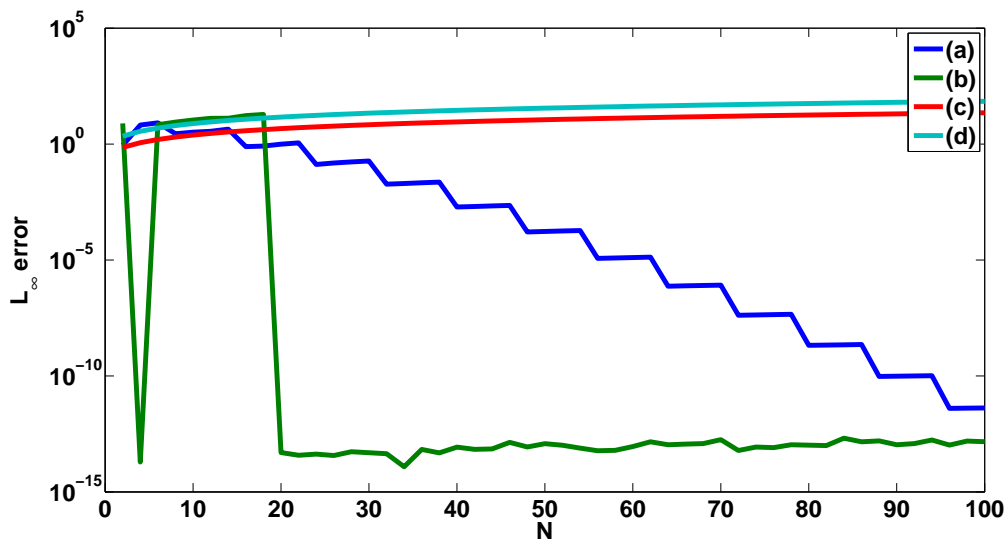
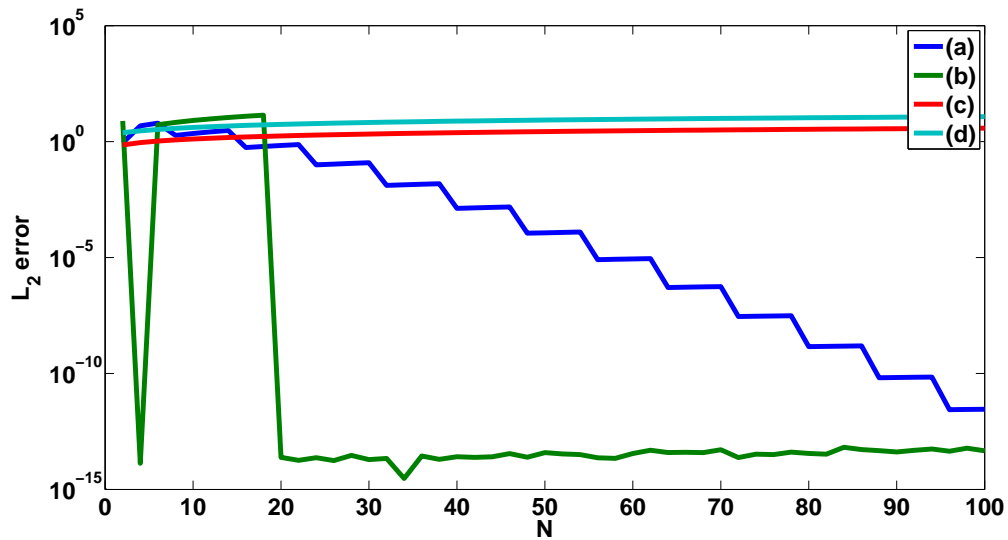
- (b)  $f(x) = \cos(10x)$
- (c)  $f(x) = \cos(x/2)$
- (d)  $f(x) = x$

All functions are defined on  $[0, 2\pi]$ .

Compute the pointwise error ( $L_\infty$ ) and the global error ( $L_2$ ), for increasing values of  $N$  and discuss the different behaviors and convergence rates based on what you know from a theoretical point of view.

**Solution:**

The figures below were generated using MatLab's `semilogy` plotting utility.



(a) First note that  $f(x)$  and all its derivatives are smoothly periodic on  $[0, 2\pi]$ , and so we expect spectral convergence in each norm. The analytical derivative is

$$\frac{du}{dx} = -4 \sin(4x) \exp(\cos(4x))$$

This can be rewritten as

$$-\frac{4}{2i} (e^{4ix} - e^{-4ix}) \sum_n (e^{4ix} + e^{-4ix})^n \frac{1}{n! 2^n}$$

Note that this expansion contains only wavenumbers  $4n$ , where  $n$  is an integer. Numerically we observe spectral convergence in both norms, with the error decreasing as we are able to resolve each of the  $4n$  wavenumbers. This corresponds to every 8 increments of  $N$ .

(b) The function  $f(x)$  and all its first derivative are smoothly periodic with wavenumber 10. Therefore the function is only resolved for  $N/2 \geq 10 \Rightarrow N \geq 20$ . Numerically we see this as the error instantly hits machine precision for  $N = 20$ , completely resolving the wave.

(c) and (d) : These functions are not periodic and therefore we do not observe convergence in either  $L_2$  or  $L_\infty$